# Sets of Marginals and Pearson-Correlation-based CHSH Inequalities for a Two-Qubit System

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> ISIT 2021 Melbourne, Australia

### **Outline**

- Introduction
  - Probability Mass Functions (PMFs)
  - Quantum Mass Functions (QMFs)
  - Simple Quantum Mass Functions (SQMFs)
- ► Main Results



### **PMFs**

Consider a sequence  $Y_1, \ldots, Y_n$  of random variables with the joint PMF

$$P_{Y_1,\ldots,Y_n}(y_1,\ldots,y_n), \qquad y_1\in\mathcal{Y}_1,\ldots,y_n\in\mathcal{Y}_n.$$

In a typical scenario of interest, we might have observed

$$Y_1 = y_1, \dots, Y_{n-1} = y_{n-1}$$

and would like to estimate  $Y_n$  based on these observations.

Usually,  $P_{Y_1,...,Y_n}(y_1,...,y_n)$  does not have a "nice" factorization.

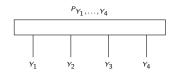
However, very often it is possible to find a function p(x, y) such that

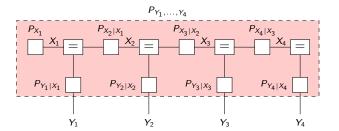
- 1.  $p(x, y) \in \mathbb{R}_{\geq 0}$  for all x, y;
- 2.  $\sum_{x,y} p(x,y) = 1$ ;
- 3.  $\sum_{\mathbf{x}} p(\mathbf{x}, \mathbf{y}) = P_{\mathbf{Y}}(\mathbf{y})$  for all  $\mathbf{y}$ ;
- **4.** p(x, y) has a "nice" factorization.

Note that p(x, y) represents a joint PMF over x and y.

### **PMFs**

### **Example (A Hidden Markov Chain)**





After applying a closing-the-box (CTB) operation to the above factor graph, i.e., summing over the variables associated with the internal edges, we obtain  $P_{Y_1,...,Y_4}$ .

## **QMFs**

Consider again a sequence  $Y_1, \ldots, Y_n$  of random variables with the joint PMF

$$P_{Y_1,\ldots,Y_n}(y_1,\ldots,y_n), \qquad y_1\in\mathcal{Y}_1,\,\ldots,\,y_n\in\mathcal{Y}_n.$$

However, now we assume that these random variables represent the measurements obtained by running some quantum-mechanical experiment.

Again, a typical scenario of interest is that we would like to estimate  $Y_n$  based on the observations

$$Y_1 = y_1, \ldots, Y_{n-1} = y_{n-1}.$$

## **QMFs**

In general, the PMF  $P_Y(y)$  does not have a "nice" factorization.

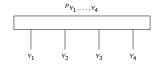
However, frequently it is possible to introduce suitable auxiliary variables  $x_1, \ldots x_m, x'_1, \ldots, x'_m$  such that there is a function  $q(\mathbf{x}, \mathbf{x}', \mathbf{y})$  satisfying

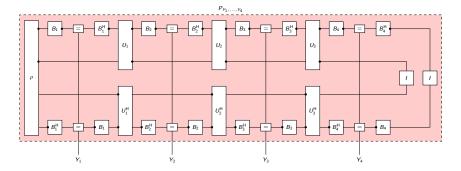
- 1.  $q(x, x', y) \in \mathbb{C}$  for all x, x', y;
- 2.  $\sum_{x,x',y} q(x,x',y) = 1;$
- 3. q(x, x', y) is a positive semi-definite (PSD) kernel in (x, x') for every y;
- 4.  $\sum_{x,x'} q(x,x',y) = P_Y(y);$
- **5.** q(x, x', y) has a "nice" factorization.

The function q is called a quantum mass function (QMF) in [Loeliger and Vontobel, 2017].

## **QMFs**

### **Example**





After applying a CTB operation to the above factor graph, i.e., summing over the variables associated with the internal edges, we obtain  $P_{Y_1,...,Y_4}$ .

## **SQMFs**

In [Loeliger and Vontobel, 2020], the authors also introduced simple quantum mass functions (SQMFs).

An SQMF q(x, x') satisfies

- 1.  $q(\mathbf{x}, \mathbf{x}') \in \mathbb{C}$  for all  $\mathbf{x}, \mathbf{x}'$ ;
- 2.  $\sum_{x,x'} q(x,x') = 1;$
- 3. q(x, x') is a PSD kernel in (x, x').

### Remark

Observations **y** in QMFs do not appear explicitly in SQMFs anymore. However, as we will see later, observations **y** emerge from SQMFs.

## **SQMFs**

### **Definition**

For  $\mathbf{x} = (x_1, \dots, x_m)$ , let  $\mathcal{I} \subseteq \{1, \dots, m\}$  and let  $\mathcal{I}^c$  be its complement.

The variables  $x_{\mathcal{I}}$  are called jointly **classicable** if the marginalized SQMF

$$q_{\mathcal{I}}(\mathbf{x}_{\mathcal{I}},\mathbf{x}_{\mathcal{I}}') := \sum_{\mathbf{x}_{\mathcal{I}^{\mathrm{c}}},\mathbf{x}_{\mathcal{I}^{\mathrm{c}}}'} q(\mathbf{x},\mathbf{x}')$$

is zero for all  $(\mathbf{x}_{\mathcal{I}}, \mathbf{x}_{\mathcal{I}}')$  satisfying  $\mathbf{x}_{\mathcal{I}} \neq \mathbf{x}_{\mathcal{I}}'$ .

### **Definition**

If the variables  $\mathbf{x}_{\mathcal{I}}$  are jointly classicable then

$$p(\mathbf{x}_{\mathcal{I}}) := q_{\mathcal{I}}(\mathbf{x}_{\mathcal{I}}, \mathbf{x}_{\mathcal{I}}), \qquad \mathbf{x}_{\mathcal{I}} \in \mathcal{X}_{\mathcal{I}},$$

represents a joint PMF over  $x_{\mathcal{I}}$ .

### **Definition**

Let  $\mathcal K$  be a collection of subsets  $\mathcal I$  of  $\{1,\ldots,m\}$  such that  $\mathbf x_{\mathcal I}$  is classicable.

## SQMFs vs. QMFs vs. PMFs

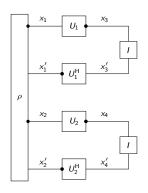
### Remark

- ▶ By defining  $p(\mathbf{x}_{\mathcal{I}}) := q_{\mathcal{I}}(\mathbf{x}_{\mathcal{I}}, \mathbf{x}_{\mathcal{I}})$ , we can see that the observations  $\mathbf{y}$  that were omitted when going from QMFs to SQMFs can "emerge" again.
- ▶ Typically, the set of marginals  $\{p(\mathbf{x}_{\mathcal{I}})\}_{\mathcal{I} \in \mathcal{K}}$  is "incompatible", i.e., there is no PMF  $p(\mathbf{x})$  such that  $p_{\mathcal{I}}(\mathbf{x}_{\mathcal{I}})$  is a marginal of  $p(\mathbf{x})$  for all  $\mathcal{I} \in \mathcal{K}$ .
- Note that there is a strong connection of SQMFs to the so-called decoherence functional [Gell-Mann and Hartle, 1989, Dowker and Halliwell, 1992], and via this also to the consistent-histories approach to quantum mechanics [Griffiths, 2002]. However, the starting point of our investigations is quite different.

### **Example**

Consider the following quantum factor graph (Q-FG), where

$$\begin{split} \mathcal{X}_i &= \mathcal{X}_i' := \{0,1\}, \qquad i \in \{1,\dots,4\}, \qquad \textit{U}_1 = \textit{U}_2 := \frac{\sqrt{2}}{2} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}, \\ \boldsymbol{\psi} &:= \begin{pmatrix} 1 & 1 & 0 \end{pmatrix}^\mathsf{T}, \qquad \boldsymbol{\rho} := \boldsymbol{\psi} \cdot \boldsymbol{\psi}^\mathsf{H}. \end{split}$$



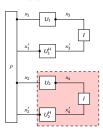
The following matrix shows the components of the SQMF  $q(\mathbf{x}, \mathbf{x}')$ , where both the row index  $(x_1, \ldots, x_4)$  and column index  $(x_1', \ldots, x_4')$  range over  $(0, 0, 0, 0), (0, 0, 0, 1), \ldots, (1, 1, 1, 1)$ .

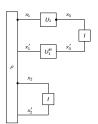
/	$\alpha_1$	0	0	0	$\alpha_1$	0	0	0	$\alpha_1$	0	0	0	0	0	0	0 \	
1	0	$\alpha_1$	0	0	0	$\beta_1$	0	0	0	$\alpha_1$	0	0	0	0	0	0	
ı	0	0	$\alpha_1$	0	0	0	$\alpha_1$	0	0	0	$\beta_1$	0	0	0	0	0	
ı	0	0	0	$\alpha_1$	0	0	0	$\beta_1$	0	0	0	$\beta_1$	0	0	0	0	
ı	$\alpha_1$	0	0	0	$\alpha_1$	0	0	0	$\alpha_1$	0	0	0	0	0	0	0	
l	0	$\beta_1$	0	0	0	$\alpha_1$	0	0	0	$\beta_1$	0	0	0	0	0	0	
ı	0	0	$\alpha_1$	0	0	0	$\alpha_1$	0	0	0	$\beta_1$	0	0	0	0	0	
ı	0	0	0	$\beta_1$	0	0	0	$\alpha_1$	0	0	0	$\alpha_1$	0	0	0	0	
ı	$\alpha_1$	0	0	0	$\alpha_1$	0	0	0	$\alpha_1$	0	0	0	0	0	0	0	
ł	0	$\alpha_1$	0	0	0	$\beta_1$	0	0	0	$\alpha_1$	0	0	0	0	0	0	
ı	0	0	$\beta_1$	0	0	0	$\beta_1$	0	0	0	$\alpha_1$	0	0	0	0	0	
ı	0	0	0	$\beta_1$	0	0	0	$\alpha_1$	0	0	0	$\alpha_1$	0	0	0	0	
ı	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	
I	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	
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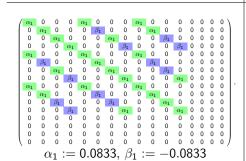
Here:  $\alpha_1 := 0.0833$ ,  $\beta_1 := -0.0833$ .

Note that the above matrix is not diagonal.

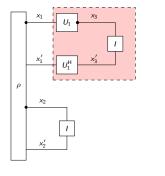
Consider  $q_{1,2,3}(x_1, x_2, x_3, x_1', x_2', x_3') := \sum_{x_4, x_4'} q(x_1, \dots, x_4, x_1', \dots, x_4')$ .

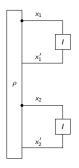






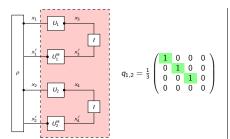
Consider  $q_{1,2}(x_1, x_2, x_1', x_2') := \sum_{x_3, x_2'} q_{1,2,3}(x_1, x_2, x_3, x_1', x_2', x_3')$ .

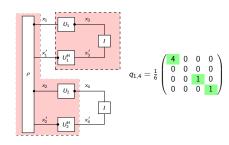


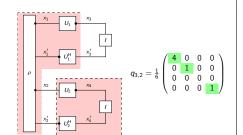


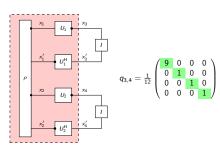
$$\frac{1}{3} \left( \begin{array}{ccccc} \mathbf{1} & 0 & 0 & 0 \\ 0 & \mathbf{1} & 0 & 0 \\ 0 & 0 & \mathbf{1} & 0 \\ 0 & 0 & 0 & 0 \end{array} \right)$$

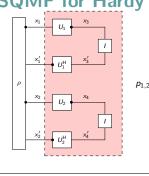
Note that the above matrix is diagonal.

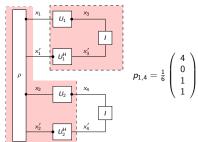


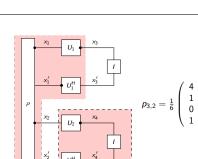


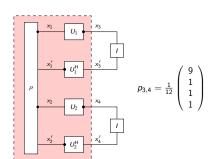












- The marginal  $p_{3,4}(1,1) = \frac{1}{12}$  shows that it is possible to have  $x_3 = x_4 = 1$ .
- The marginals  $p_{3,2}(1,0) = 0$  and  $p_{3,2}(1,1) = 1/6$  show that the condition  $x_3 = 1$  implies  $x_2 = 1$ .
- The marginals  $p_{1,4}(0,1) = 0$  and  $p_{1,4}(1,1) = 1/6$  show that the condition  $x_4 = 1$  implies  $x_1 = 1$ .
- However, the marginal  $p_{1,2}(1,1) = 0$  implies that we cannot have  $x_1 = x_2 = 1$ , which contradicts  $p_{3,4}(1,1) > 0$ .

### Remark

We have expressed Hardy's paradox in terms of marginals of SQMFs.

Other paradoxes (e.g. Bell's test, Wigner's friend experiment, and the Frauchiger-Renner paradox) can also be expressed in terms of some suitably defined SQMFs.

## **SQMFs**

### Remark

▶ For any two sets  $\mathcal{I}_1$ ,  $\mathcal{I}_2 \in \mathcal{K}$ , the following local consistency constraint holds:

$$\sum_{m{x}_{\mathcal{I}_1 \setminus \mathcal{I}_2}} p(m{x}_{\mathcal{I}_1}) = \sum_{m{x}_{\mathcal{I}_2 \setminus \mathcal{I}_1}} p(m{x}_{\mathcal{I}_2}) \quad ext{(for all } m{x}_{\mathcal{I}_1 \cap \mathcal{I}_2}).$$

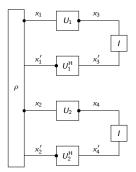
▶ Interestingly, this requirement is very similar to the properties of the beliefs in the local marginal polytope (LMP) of a standard factor graph (S-FG) [Wainwright and Jordan, 2008].

## **SQMFs**

There are two extreme cases to be considered:

- **1.** The set of marginals  $\{p_{\mathcal{I}}\}_{\mathcal{I}\in\mathcal{K}}$  can be achieved by some joint PMF.
- 2. The set of marginals  $\{p_{\mathcal{I}}\}_{\mathcal{I} \in \mathcal{K}}$  satisfies only the local consistency constraints, i.e.,  $\{p_{\mathcal{I}}\}_{\mathcal{I} \in \mathcal{K}}$  is in the LMP.

Consider the set of marginals created by jointly classicable variables in the following quantum system.



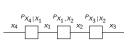
How general can the marginals  $\{p_{\mathcal{I}}\}_{\mathcal{I}\in\mathcal{K}}$  be for different  $\rho$ ,  $U_1$ , and  $U_2$ ?

### **Definition**

### We define

- 1.  $\mathcal{M}(N)$  to be the set of realizable marginals of the S-FG  $N \in \{N_1, N_2, N_3\}$ , where the local functions in N are varied;
- 2.  $\mathcal{LM}(\mathcal{K})$  to be the LMP of the S-FG  $N_1$ ;
- 3.  $\mathcal{M}(N_4)$  to be the set of the classicable variables' marginals in the two-qubit system  $N_4$ , where  $\rho$ ,  $U_1$ , and  $U_2$  are varied.







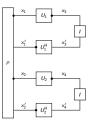


Figure: The S-FG  $N_1$ .

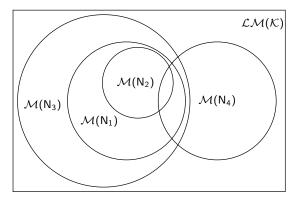
Figure: The S-FG N<sub>2</sub>.

Figure: The S-FG N<sub>3</sub>.

Figure: The Q-FG N<sub>4</sub>.

### **Theorem**

The following Venn diagram holds.



We prove that each part in the diagram is non-empty.

Let us consider the random variables  $X_1,\ldots,X_4\in\{0,1\}$ . The Clauser-Horne-Shimony-Holt (CHSH) inequality states that

$$\left|\mathbb{E}(X_1\cdot X_2)+\mathbb{E}(X_1\cdot X_4)+\mathbb{E}(X_3\cdot X_2)-\mathbb{E}(X_3\cdot X_4)\right|\leq 2.$$

In this paper, we prove a Pearson correlation coefficient (PCC)-based variant of the CHSH inequality.

### **Theorem**

Suppose that the random variables  $X_1, \ldots, X_4 \in \{0,1\}$  satisfy

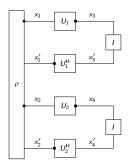
$$\operatorname{Var}(X_1), \ldots, \operatorname{Var}(X_4) \in \mathbb{R}_{>0}.$$

Then the following PCC-based CHSH inequality holds:

$$\left|\operatorname{Corr}(X_1 \cdot X_2) + \operatorname{Corr}(X_1 \cdot X_4) + \operatorname{Corr}(X_3 \cdot X_2) - \operatorname{Corr}(X_3 \cdot X_4)\right| \leq \frac{5}{2}.$$

This resolves a conjecture proposed in [Pozsgay et al., 2017].

### **Conclusion**



### **Proposition**

- ► The SQMF of the above Q-FG can lead to "incompatible" marginals.
- We characterize the relationships among the sets of marginals mentioned in the previous slides.
- Many well-known quantum phenomena, e.g., Hardy's paradox and Bell's test, can be cast with this SQMF.

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# Thank you!

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